

## Differential equations at resonance

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*Abstract.* New existence results are presented for the two point singular “resonant” boundary value problem  $\frac{1}{p}(py')' + ry + \lambda_m qy = f(t, y, py')$  a.e. on  $[0, 1]$  with  $y$  satisfying Sturm Liouville or Periodic boundary conditions. Here  $\lambda_m$  is the  $(m + 1)^{st}$  eigenvalue of  $\frac{1}{pq}[(pu')' + rpu] + \lambda u = 0$  a.e. on  $[0, 1]$  with  $u$  satisfying Sturm Liouville or Periodic boundary data.

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### 1. Introduction

In this paper we derive some existence results for the second order equation

$$(1.1) \quad \frac{1}{p(t)}(p(t)y'(t))' + r(t)y(t) + \lambda_m q(t)y(t) = f(t, y(t), p(t)y'(t)) \quad \text{a.e. on } [0, 1]$$

with  $y$  satisfying either

(i) (Sturm Liouville)

$$(SL) \quad \begin{cases} -\alpha y(0) + \beta \lim_{t \rightarrow 0^+} p(t)y'(t) = 0, & \alpha \geq 0, \beta \geq 0, \alpha^2 + \beta^2 > 0 \\ \alpha y(1) + b \lim_{t \rightarrow 1^-} p(t)y'(t) = 0, & a \geq 0, b \geq 0, a^2 + b^2 > 0 \end{cases}$$

or

(ii) (Periodic)

$$(P) \quad \begin{cases} y(0) = y(1) \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = \lim_{t \rightarrow 1^-} p(t)y'(t). \end{cases}$$

*Remarks.* (i)  $\lambda_m$  will be described later.

(ii) The Neumann condition  $\lim_{t \rightarrow 0^+} p(t)y'(t) = \lim_{t \rightarrow 1^-} p(t)y'(t) = 0$  is included in (SL) with  $\alpha = a = 0$ .

(iii) If a function  $u \in C[0, 1] \cap C^1(0, 1)$  with  $pu' \in C[0, 1]$  satisfies boundary condition (i) we write  $u \in (SL)$ . A similar remark applies for the boundary condition (ii).

Throughout the paper  $p \in C[0, 1] \cap C^1(0, 1)$  together with  $p > 0$  on  $(0, 1)$ . Also  $pf : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$  is an  $L^1$ -Carathéodory function. By this we mean:

- (i)  $t \rightarrow p(t)f(t, y, q)$  is measurable for all  $(y, q) \in \mathbf{R}^2$ ;
- (ii)  $(y, q) \rightarrow p(t)f(t, y, q)$  is continuous for a.e.  $t \in [0, 1]$ ;

- (iii) for any  $r > 0$  there exists  $h_r \in L^1[0, 1]$  such that  $|p(t)f(t, y, q)| \leq h_r(t)$  for a.e.  $t \in [0, 1]$  and for all  $|y| \leq r, |q| \leq r$ .

For notational purposes let  $w$  be a weight function. By  $L^1_w[0, 1]$  we mean the space of functions  $u$  such that  $\int_0^1 w(t)|u(t)| dt < \infty$ .  $L^2_w[0, 1]$  denotes the space of functions  $u$  such that  $\int_0^1 w(t)|u(t)|^2 dt < \infty$ ; also for  $u, v \in L^2_w[0, 1]$  define  $\langle u, v \rangle = \int_0^1 w(t)u(t)\overline{v(t)} dt$ . Let  $AC[0, 1]$  be the space of functions which are absolutely continuous on  $[0, 1]$ .

Before we discuss the boundary value problem (1.1) and its appropriate literature we first gather together some facts on second order differential equations ([12], [16]). Consider the linear equation

$$(1.2) \quad \begin{cases} \frac{1}{p}(py')' + \tau y = g(t) & \text{a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}). \end{cases}$$

By a solution to (1.2) we mean a function  $y \in C[0, 1] \cap C^1(0, 1)$  with  $py' \in AC[0, 1]$  which satisfies the differential equation in (1.2) a.e. on  $[0, 1]$  and the stated boundary conditions.

**Theorem 1.1.** *Suppose*

$$(1.3) \quad p \in C[0, 1] \cap C^1(0, 1) \text{ with } p > 0 \text{ on } (0, 1) \text{ and } \int_0^1 \frac{ds}{p(s)} < \infty$$

and

$$(1.4) \quad \tau, g \in L^1_p[0, 1]$$

are satisfied. If

$$(1.5) \quad \begin{cases} \frac{1}{p}(py')' + \tau y = 0 & \text{a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

has only the trivial solution, then (1.2) has exactly one solution  $y$  given by

$$y(t) = d_0u_1(t) + d_1u_2(t) + \int_0^t \frac{[u_2(t)u_1(s) - u_1(t)u_2(s)]}{W(s)}g(s) ds$$

where  $u_1$  is the unique solution to

$$\begin{cases} \frac{1}{p}(pu')' + \tau u = 0 & \text{a.e. on } [0, 1] \\ u(0) = 1, \lim_{t \rightarrow 0^+} p(t)u'(t) = 0 \end{cases}$$

and  $u_2$  is the unique solution to

$$\begin{cases} \frac{1}{p}(pu')' + \tau u = 0 & \text{a.e. on } [0, 1] \\ u(0) = 0, \lim_{t \rightarrow 0^+} p(t)u'(t) = 1 \end{cases}$$

and  $d_0$  and  $d_1$  are uniquely determined from the boundary condition;  $W$  of course denotes the Wronskian. In fact

$$y(t) = \int_0^1 G(t, s)g(s) ds$$

with

$$G(t, s) = \begin{cases} \frac{y_1(s)y_2(t)}{W(s)}, & 0 < s \leq t \\ \frac{y_1(t)y_2(s)}{W(s)}, & t \leq s < 1 \end{cases}$$

where  $y_1$  and  $y_2$  are the two “usual” linearly independent solutions i.e. choose  $y_1 \neq 0, y_2 \neq 0$  so that  $y_1, y_2$  satisfy  $\frac{1}{p}(py')' + \tau y = 0$  a.e. on  $[0, 1]$  with  $y_1$  satisfying the first boundary condition and  $y_2$  satisfying the second boundary condition.

We now state an existence principle ([16]), which was established using fixed point methods, for the second order nonresonant boundary value problem

$$(1.6) \quad \begin{cases} \frac{1}{p}(py')' + \tau y = f(t, y, py') & \text{a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}). \end{cases}$$

**Theorem 1.2.** Let  $pf : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$  be an  $L^1$ -Carathéodory function and assume (1.3) and

$$(1.7) \quad \tau \in L^1_p[0, 1]$$

hold. In addition suppose (1.5) has only the trivial solution. Now assume there is a constant  $M_0$ , independent of  $\lambda$ , with

$$\|y\|_* = \max\{\sup_{[0,1]} |y(t)|, \sup_{(0,1)} |p(t)y'(t)|\} \leq M_0$$

for any solution  $y$  to

$$\begin{cases} \frac{1}{p}(py')' + \tau y = \lambda f(t, y, py') & \text{a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

for each  $\lambda \in (0, 1)$ . Then (1.6) has at least one solution  $u \in C[0, 1] \cap C^1(0, 1)$  with  $pu' \in AC[0, 1]$ .

Next we gather together some results on the Sturm Liouville eigenvalue problem

$$(1.8) \quad \begin{cases} Lu = \lambda u & \text{a.e. on } [0, 1] \\ u \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

where  $Lu = -\frac{1}{pq(t)}[(pu')' + r(t)pu]$ . Assume (1.3) and

$$(1.9) \quad r, q \in L^1_p[0, 1] \text{ with } q > 0 \text{ a.e. on } [0, 1]$$

hold. Let

$$D(L) = \{w \in C[0, 1] : w, pw' \in AC[0, 1] \text{ with } w \in (\text{SL}) \text{ or } (\text{P})\}.$$

Then  $L$  has a countably infinite number ([1], [12], [16]) of real eigenvalues  $\lambda_i$  with corresponding eigenfunctions  $\psi_i \in D(L)$ . The eigenfunctions  $\psi_i$  may be chosen so that they form an orthonormal set and we may also arrange the eigenvalues so that

$$(1.10) \quad \lambda_0 < \lambda_1 < \lambda_2 < \dots$$

*Remark.* The  $\lambda_i$ 's may be estimated numerically ([2]) using SLEIGN.

In addition the set of eigenfunctions  $\psi_i$  form a basis for  $L^2_{pq}[0, 1]$  and if  $h \in L^2_{pq}[0, 1]$  then  $h$  has a Fourier series representation and  $h$  satisfies Parseval's equality i.e.

$$h = \sum_{i=0}^{\infty} \langle h, \psi_i \rangle \psi_i \quad \text{and} \quad \int_0^1 pq|h|^2 dt = \sum_{i=0}^{\infty} |\langle h, \psi_i \rangle|^2.$$

We are concerned with existence results for the nonlinear second order equation

$$(1.11) \quad \begin{cases} \frac{1}{p}(py')' + ry + \lambda_m qy = f(t, y, py') & \text{a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

where  $\lambda_m$  is the  $(m + 1)^{st}$  eigenvalue of (1.8). In recent years several authors ([4], [7]–[9], [11], [13], [18]–[19]) have examined the boundary value problems

$$\begin{cases} y'' + n^2\pi^2y = f(t, y) & \text{a.e. on } [0, 1] \\ y(0) = y(1) = 0 \end{cases}$$

and

$$\begin{cases} y'' + m^2\pi^2y = f(t, y) & \text{a.e. on } [0, 1] \\ y(0) = y(1), y'(0) = y'(1) \end{cases}$$

where  $n \geq 1, m \geq 0$  are integers. Most of the papers in the literature ([3], [7], [11], [18]–[19]) concentrate on the first eigenvalue ( $n = 1$  or  $m = 0$ ). However over the last ten years or so ([6], [10]) the case when  $n > 1$  or  $m > 0$  has been discussed. This paper continues this study for the more general problem (1.11); also it provides a new approach to studying the above resonant type problems. We refer the reader to [6]–[9] for many of the motivating ideas in this paper. Finally it is of interest to note that in previous studies ([6], [8], [11]) the nonlinearity  $f$  is required to grow no more than linearly in  $y$  as  $|y| \rightarrow \infty$  whereas in this paper solutions will exist provided  $f$  grows fast enough e.g.  $yf(t, y, z) \geq A|y|^{\theta+1}$  for some  $A > 0$  and  $\theta > 0$ .

**2. Existence**

Existence theory is developed for the second order boundary value problem

$$(2.1) \quad \begin{cases} \frac{1}{p}(py')' + ry + \lambda_m qy = f(t, y, py') & \text{a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

where  $\lambda_m$  is the  $(m + 1)^{\text{st}}$  eigenvalue of

$$(2.2) \quad \begin{cases} Lu = \lambda u & \text{a.e. on } [0, 1] \\ u \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

and  $Lu = -\frac{1}{pq(t)}[(pu')' + r(t)pu]$ .

Two types of existence results are presented, the first examines the problem on the “left” of the eigenvalue whereas the second discusses the problem on the “right” of the eigenvalue.

**Existence theory I.**

Throughout this subsection let

$$H_{\alpha_0, \theta}(u_1) = \begin{cases} |u_1|^{\theta+1}, & |u_1| \leq 1 \\ |u_1|^{\alpha_0+1}, & |u_1| > 1. \end{cases}$$

**Theorem 2.1.** *Let  $pf : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$  be an  $L^1$ -Carathéodory function with (1.3) and (1.9) satisfied. Suppose  $f$  has the decomposition  $f(t, u_1, u_2) = g(t, u_1, u_2) + h(t, u_1, u_2)$  with  $pg, ph : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$   $L^1$ -Carathéodory functions and*

$$(2.3) \quad \begin{cases} \text{there exist constants } A > 0, 0 < \alpha_0 < 1 \text{ and a function} \\ \phi \in L^1_p[0, 1], \phi > 0 \text{ a.e. on } [0, 1] \text{ with } u_1 g(t, u_1, u_2) \geq A\phi(t)H_{\alpha_0, \theta}(u_1) \\ \text{for a.e. } t \in [0, 1]; \text{ here } \alpha_0 \leq \theta \end{cases}$$

$$(2.4) \quad \left\{ \begin{array}{l} \text{there exist } \phi_i \in L_p^1[0, 1], i = 1, 2, 3 \text{ and constants } \beta_0 \text{ and} \\ \sigma \text{ with } |h(t, u_1, u_2)| \leq \phi_1(t) + \phi_2(t)|u_1|^{\beta_0} + \phi_3(t)|u_2|^\sigma \text{ for} \\ \text{a.e. } t \in [0, 1]; \text{ here } \beta_0 < \alpha_0 \text{ and } 0 \leq \sigma < \frac{\alpha_0}{2} \text{ and} \\ \phi_3 > 0 \text{ a.e. on } [0, 1] \text{ or } \phi_3 \equiv 0 \text{ on } [0, 1] \end{array} \right.$$

$$(2.5) \quad \left\{ \begin{array}{l} \text{there exist } \phi_i \in L_p^1[0, 1], i = 4, 5 \text{ and a constant } \gamma \leq \alpha_0 \text{ with} \\ |g(t, u_1, u_2)| \leq \phi_4(t) + \phi_5(t)|u_1|^\gamma \text{ for a.e. } t \in [0, 1] \end{array} \right.$$

$$(2.6) \quad \left\{ \begin{array}{l} \phi_4^2 q^{-1} \in L_p^1[0, 1], \phi_1^2 q^{-1} \in L_p^1[0, 1], \\ \left( \phi_5^{2(\alpha_0+1)} q^{-(\alpha_0+1)} \phi^{-2\gamma} \right)^{\frac{1}{\alpha_0+1-2\gamma}} \in L_p^1[0, 1], \\ \left( \phi_2^{2(\alpha_0+1)} q^{-(\alpha_0+1)} \phi^{-2\beta_0} \right)^{\frac{1}{\alpha_0+1-2\beta_0}} \in L_p^1[0, 1] \text{ and} \\ \left( \phi^2 q^{-(\alpha_0+1)} \right)^{\frac{1}{1-\alpha_0}} \in L_p^1[0, 1] \end{array} \right.$$

and

$$(2.7) \quad \left\{ \begin{array}{l} \left( \phi_1^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} \in L_p^1[0, 1], \left( \phi_2^{\alpha_0+1} \phi^{-(\beta_0+1)} \right)^{\frac{1}{\alpha_0-\beta_0}} \in L_p^1[0, 1], \\ \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} \in L_p^1[0, 1], \\ \left( \phi_5^{\alpha_0+1} \phi^{-\gamma} \right)^{\frac{1}{\alpha_0+1-\gamma}} \in L_p^1[0, 1] \text{ and} \\ \left( q^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} \in L_p^1[0, 1] \end{array} \right.$$

holding. Then (2.1) has at least one solution  $y \in C[0, 1] \cap C^1(0, 1)$  with  $py' \in AC[0, 1]$ .

*Remark.* Typical examples where (2.3) is satisfied are say (i)  $g(t, u_1, u_2) = u_1^{\frac{m}{n}}$ ,  $m$  odd and  $n$  odd or (ii)  $g(t, u_1, u_2) = u_1^{\frac{1}{2}}$ ,  $u_1 \geq 0$  with  $g(t, u_1, u_2) = -|u_1|^{\frac{1}{2}}$ ,  $u_1 < 0$ .

PROOF: Consider the family of problems

$$(2.8)_\lambda \quad \left\{ \begin{array}{l} \frac{1}{p}(py')' + ry + \mu qy = \lambda[f(t, y, py') + (\mu - \lambda_m)qy] \text{ a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}) \end{array} \right.$$

where  $0 < \lambda < 1$  and  $\lambda_{m-1} < \mu < \lambda_m$ ; here  $\lambda_{-1} = -\infty$  (for notational purposes) with  $\lambda_i$  as described in (1.10).

Notice  $L_{pq}^2[0, 1] = \Omega \oplus \Omega^\perp$  where  $\Omega = span\{\psi_0, \psi_1, \dots, \psi_{m-1}\}$ ; here  $\psi_i$  are the eigenfunctions corresponding to the eigenvalues  $\lambda_i$  (see Section 1).

Let  $y$  be any solution to  $(2.8)_\lambda$ . Then  $y = u + w$  where  $u \in \Omega$  and  $w \in \Omega^\perp$ . Multiply  $(2.8)_\lambda$  by  $w - u$  and integrate from 0 to 1 to obtain

$$\begin{aligned} \int_0^1 (w - u)(py')' dt + \int_0^1 pr[w^2 - u^2] dt + \mu \int_0^1 pq[w^2 - u^2] dt \\ = \lambda \int_0^1 (w - u)pf(t, y, py') dt + \lambda(\mu - \lambda_m) \int_0^1 pq[w^2 - u^2] dt. \end{aligned}$$

Integration by parts yields

$$\int_0^1 (w - u)(py')' dt = Q_0 - \int_0^1 p(w')^2 dt + \int_0^1 p(u')^2 dt$$

where

$$Q_0 = \begin{cases} -\frac{a}{b}[w^2(1) - u^2(1)] - \frac{a}{\beta}[w^2(0) - u^2(0)] & \text{if } y \in (SL) \\ 0 & \text{if } y \in (P); \end{cases}$$

here  $y(0) = 0$  means  $u(0) + w(0) = 0$  and so  $u(0) = w(0) = 0$ . Thus we have

$$\begin{aligned} (2.9) \quad Q_0 + \int_0^1 [-p(w')^2 + prw^2 + \mu pqw^2] dt + \int_0^1 [p(u')^2 - pru^2 - \mu pq u^2] dt \\ = \lambda \int_0^1 (w - u)pf(t, y, py') dt + \lambda(\mu - \lambda_m) \int_0^1 pqw^2 dt \\ - \lambda(\mu - \lambda_m) \int_0^1 pq u^2 dt. \end{aligned}$$

Now since  $u \in \Omega$ ,  $w \in \Omega^\perp$  and  $y = u + w$  we have

$$u = \sum_{i=0}^{m-1} c_i \psi_i \quad \text{and} \quad w = \sum_{i=m}^{\infty} c_i \psi_i \quad \text{where} \quad c_i = \langle y, \psi_i \rangle;$$

note  $u = 0$  if  $m = 0$ . Then since  $(p\psi'_i)' + r p\psi_i + \lambda_i p q \psi_i = 0$  we have

$$\begin{aligned} Q_0 + \int_0^1 [-p(w')^2 + prw^2 + \mu pqw^2] dt + \int_0^1 [p(u')^2 - pru^2 - \mu pq u^2] dt \\ = \sum_{i=m}^{\infty} (\mu - \lambda_i) c_i^2 \int_0^1 pq \psi_i^2 dt + \sum_{i=0}^{m-1} (\lambda_i - \mu) c_i^2 \int_0^1 pq \psi_i^2 dt \\ \leq (\mu - \lambda_m) \int_0^1 pqw^2 dt + (\lambda_{m-1} - \mu) \int_0^1 pq u^2 dt. \end{aligned}$$

Put this into (2.9) to obtain

$$\begin{aligned} & \lambda \int_0^1 (w - u)pg(t, y, py') dt + (1 - \lambda)(\lambda_m - \mu) \int_0^1 pqw^2 dt \\ & + (\mu - \lambda_{m-1}) \int_0^1 pqu^2 dt + \lambda(\lambda_m - \mu) \int_0^1 pqu^2 dt \\ & \leq -\lambda \int_0^1 (w - u)ph(t, y, py') dt. \end{aligned}$$

Consequently

$$(2.10) \quad \begin{aligned} & \int_0^1 pyg(t, y, py') dt + (\lambda_m - \mu) \int_0^1 pqu^2 dt \leq 2 \int_0^1 pug(t, y, py') dt \\ & + \int_0^1 p|y||h(t, y, py')| dt + 2 \int_0^1 p|u||h(t, y, py')| dt. \end{aligned}$$

Assumption (2.3) yields

$$\begin{aligned} \int_0^1 pyg(t, y, py') dt & \geq A \int_0^1 p\phi H_{\alpha_0, \theta}(y) dt \\ & = A \int_0^1 p\phi|y|^{\alpha_0+1} dt + A \int_{\{t:|y(t)| \leq 1\}} p\phi[|y|^{\theta+1} - |y|^{\alpha_0+1}] dt \\ & \geq A \int_0^1 p\phi|y|^{\alpha_0+1} dt - A \int_0^1 p\phi dt \end{aligned}$$

and put this into (2.10), and use (2.4) and (2.5), to obtain

$$(2.11) \quad \begin{aligned} & A \int_0^1 p\phi|y|^{\alpha_0+1} dt + (\lambda_m - \mu) \int_0^1 pqu^2 dt \leq A \int_0^1 p\phi dt + 2 \int_0^1 p\phi_4|u| dt \\ & + 2 \int_0^1 p\phi_5|u||y|^\gamma dt + \int_0^1 p\phi_1|y| dt \\ & + \int_0^1 p\phi_2|y|^{\beta_0+1} dt + \int_0^1 p\phi_3|y||py'|^\sigma dt \\ & + 2 \int_0^1 p\phi_1|u| dt + 2 \int_0^1 p\phi_2|u||y|^{\beta_0} dt \\ & + 2 \int_0^1 p\phi_3|u||py'|^\sigma dt. \end{aligned}$$

For the remainder of the proof we assume without loss of generality that  $\sigma > 0$  and  $\phi_3 \neq 0$  on  $[0, 1]$ . Let  $\epsilon > 0$  be given. Hölder's inequality together with assumption (2.6) immediately yields the following inequalities:

$$2 \int_0^1 p\phi_4|u| dt \leq 2Q_1 \left( \int_0^1 pqu^2 dt \right)^{\frac{1}{2}} \leq \epsilon \int_0^1 pqu^2 dt + \frac{Q_1}{\epsilon};$$



$$\begin{aligned}
 2 \int_0^1 p\phi_1|u| dt &\leq \epsilon \int_0^1 pqu^2 dt + \frac{Q_2}{\epsilon}; \\
 2 \int_0^1 p\phi_5|u||y|^\gamma dt &\leq 2Q_3 \left( \int_0^1 pqu^2 dt \right)^{\frac{1}{2}} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\gamma}{\alpha_0+1}} \\
 &\leq \epsilon Q_3 \int_0^1 pqu^2 dt + \frac{Q_3}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\gamma}{\alpha_0+1}}; \\
 2 \int_0^1 p\phi_2|u||y|^{\beta_0} dt &\leq \epsilon Q_4 \int_0^1 pqu^2 dt + \frac{Q_4}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\beta_0}{\alpha_0+1}}; \\
 \int_0^1 p\phi_1|y| dt &\leq Q_5 \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}}; \\
 \int_0^1 p\phi_2|y|^{\beta_0+1} dt &\leq Q_6 \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\beta_0+1}{\alpha_0+1}}; \\
 \int_0^1 p\phi_3|y||py'|^\sigma dt &\leq \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}} \\
 &\quad \times \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{\alpha_0}{\alpha_0+1}}; \\
 2 \int_0^1 p\phi_3|u||py'|^\sigma dt &\leq 2Q_7 \left( \int_0^1 pqu^2 dt \right)^{\frac{1}{2}} \\
 &\quad \times \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{\alpha_0}{\alpha_0+1}} \\
 &\leq \epsilon Q_7 \int_0^1 pqu^2 dt \\
 &\quad + \frac{Q_7}{\epsilon} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{2\alpha_0}{\alpha_0+1}}
 \end{aligned}$$

for some constants  $Q_1, \dots, Q_7$ . Put these into (2.11) to obtain

$$\begin{aligned}
 &A \int_0^1 p\phi|y|^{\alpha_0+1} dt + (\lambda_m - \mu - 2\epsilon - \epsilon Q_3 - \epsilon Q_4 - \epsilon Q_7) \int_0^1 pqu^2 dt \\
 &\leq Q_8 + \frac{Q_3}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\gamma}{\alpha_0+1}} + \frac{Q_4}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\beta_0}{\alpha_0+1}} \\
 &+ Q_5 \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}} + Q_6 \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\beta_0+1}{\alpha_0+1}} \\
 &+ \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{\alpha_0}{\alpha_0+1}}
 \end{aligned}$$

$$+ \frac{Q_7}{\epsilon} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{2\alpha_0}{\alpha_0+1}}$$

for some constant  $Q_8$ . We may choose  $\epsilon$  so that  $\lambda_m - \mu - 2\epsilon - \epsilon Q_3 - \epsilon Q_4 - \epsilon Q_7 > 0$  and we have

$$\begin{aligned} A \int_0^1 p\phi|y|^{\alpha_0+1} dt &\leq Q_8 + \frac{Q_3}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\gamma}{\alpha_0+1}} \\ &+ \frac{Q_4}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\beta_0}{\alpha_0+1}} + Q_5 \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}} \\ (2.12) \quad &+ Q_6 \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\beta_0+1}{\alpha_0+1}} \\ &+ \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{\alpha_0}{\alpha_0+1}} \\ &+ \frac{Q_7}{\epsilon} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{2\alpha_0}{\alpha_0+1}}. \end{aligned}$$

We now consider two cases  $\int_0^1 p\phi|y|^{\alpha_0+1} dt > 1$  and  $\int_0^1 p\phi|y|^{\alpha_0+1} dt \leq 1$  separately.

Case (i).  $\int_0^1 p\phi|y|^{\alpha_0+1} dt > 1$ .

Divide (2.12) by  $\left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}}$  and use  $\int_0^1 p\phi|y|^{\alpha_0+1} dt > 1$  to obtain

$$\begin{aligned} A \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\alpha_0}{\alpha_0+1}} &\leq Q_8 + \frac{Q_3}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\gamma-1}{\alpha_0+1}} \\ &+ \frac{Q_4}{\epsilon} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{2\beta_0-1}{\alpha_0+1}} \\ &+ Q_5 + Q_6 \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\beta_0}{\alpha_0+1}} \\ &+ \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{\alpha_0}{\alpha_0+1}} \\ &+ \frac{Q_7}{\epsilon} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{2\alpha_0}{\alpha_0+1}}. \end{aligned}$$

Now since  $\max\{2\gamma - 1, 2\beta_0 - 1, \beta_0\} < \alpha_0$  there exist constants  $Q_9, Q_{10}$  and  $Q_{11}$

with

$$\begin{aligned} \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{\alpha_0}{\alpha_0+1}} &\leq Q_9 + Q_{10} \left(\int_0^1 p\left(\phi_3^{\alpha_0+1}\phi^{-1}\right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt\right)^{\frac{\alpha_0}{\alpha_0+1}} \\ &\quad + Q_{11} \left(\int_0^1 p\left(\phi_3^{\alpha_0+1}\phi^{-1}\right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt\right)^{\frac{2\alpha_0}{\alpha_0+1}}. \end{aligned}$$

Using the inequality  $(a + b)^c \leq 2^c(a^c + b^c)$  for  $a \geq 0, b \geq 0, c \geq 0$  we see that there exist constants  $Q_{12}$  and  $Q_{13}$  with

$$(2.13) \quad \int_0^1 p\phi|y|^{\alpha_0+1} dt \leq Q_{12} + Q_{13} \left(\int_0^1 p\left(\phi_3^{\alpha_0+1}\phi^{-1}\right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt\right)^2.$$

Case (ii).  $\int_0^1 p\phi|y|^{\alpha_0+1} dt \leq 1$ .

In this case (2.13) is clearly true with  $Q_{12} = 1$ .

Consequently in all cases (2.13) is true. Returning to  $(2.8)_\lambda$  we have

$$(2.14) \quad y(t) = \lambda \int_0^1 G(t, s)[f(s, y(s), p(s)y'(s)) + (\mu - \lambda_m)q(s)y(s)] ds$$

and

$$(2.15) \quad p(t)y'(t) = \lambda \int_0^1 p(t)G_t(t, s)[f(s, y(s), p(s)y'(s)) + (\mu - \lambda_m)q(s)y(s)] ds$$

where  $G(t, s)$  is the Green's function associated with  $\frac{1}{p}(pv')' + rv + \mu qv = 0$  a.e. on  $[0, 1]$ ,  $v \in (SL)$  or  $(P)$ .

Notice ([16], [17]) that  $\sup_{t \in [0,1]} |p(t)G_t(t, s)| \leq Q_{14}p(s)$  for some constant  $Q_{14}$ . Now (2.15) together with (2.4) and (2.5) imply for  $t \in (0, 1)$  that

$$\begin{aligned} |p(t)y'(t)| &\leq Q_{15} \int_0^1 p\phi_1 ds + Q_{15} \int_0^1 p\phi_2|y|^{\beta_0} ds + Q_{15} \int_0^1 p\phi_3|py'|^\sigma ds \\ &\quad + Q_{15} \int_0^1 p\phi_4 ds + Q_{15} \int_0^1 p\phi_5|y|^\gamma ds + Q_{16} \int_0^1 pq|y| ds \end{aligned}$$

for some constants  $Q_{15}$  and  $Q_{16}$ . Hölder's inequality together with (2.6) implies

$$\begin{aligned} |p(t)y'(t)| &\leq Q_{17} + Q_{18} \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{\beta_0}{\alpha_0+1}} \\ &\quad + Q_{19} \left(\int_0^1 p\left(\phi_3^{\alpha_0+1}\phi^{-1}\right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt\right)^{\frac{\alpha_0}{\alpha_0+1}} \\ &\quad + Q_{20} \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{\gamma}{\alpha_0+1}} + Q_{21} \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{1}{\alpha_0+1}} \end{aligned}$$

for some constants  $Q_{17}, \dots, Q_{21}$ . Thus for  $t \in (0, 1)$  we have

$$\begin{aligned}
 |p(t)y'(t)|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} &\leq Q_{22} + Q_{23} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\sigma\beta_0}{\alpha_0}} \\
 &+ Q_{24} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^\sigma \\
 (2.16) \quad &+ Q_{25} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\sigma\gamma}{\alpha_0}} \\
 &+ Q_{26} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\sigma}{\alpha_0}}
 \end{aligned}$$

for some constants  $Q_{22}, \dots, Q_{26}$ . This together with (2.13) implies

$$\begin{aligned}
 &\int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \\
 &\leq Q_{27} + Q_{28} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{2\sigma\beta_0}{\alpha_0}} \\
 &+ Q_{29} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^\sigma \\
 &+ Q_{30} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{2\sigma\gamma}{\alpha_0}} \\
 &+ Q_{31} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{2\sigma}{\alpha_0}}
 \end{aligned}$$

for some constants  $Q_{27}, \dots, Q_{31}$ . Finally since  $\max\{\frac{2\sigma\beta_0}{\alpha_0}, \sigma, \frac{2\sigma\gamma}{\alpha_0}, \frac{2\sigma}{\alpha_0}\} < 1$  there exists a constant  $Q_{32}$  with

$$(2.17) \quad \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \leq Q_{32}$$

and this together with (2.13) implies that there exists a constant  $Q_{33}$  with

$$(2.18) \quad \int_0^1 p\phi|y|^{\alpha_0+1} dt \leq Q_{33}.$$

Putting these inequalities into (2.16) establishes the existence of a constant  $Q_{34}$  with

$$(2.19) \quad \sup_{t \in (0,1)} |p(t)y'(t)| \leq Q_{34}.$$

Now (2.14) together ([16], [17]) with  $\sup_{t \in [0,1]} |G(t, s)| \leq Q_{35}p(s)$ , for some constant  $Q_{35}$ , and Hölder's inequality implies for  $t \in [0, 1]$  that

$$\begin{aligned}
 |y(t)| \leq & Q_{36} + Q_{37} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\beta_0}{\alpha_0+1}} \\
 & + Q_{38} \left( \int_0^1 p \left( \phi_3^{\alpha_0+1} \phi^{-1} \right)^{\frac{1}{\alpha_0}} |py'|^{\frac{\sigma(\alpha_0+1)}{\alpha_0}} dt \right)^{\frac{\alpha_0}{\alpha_0+1}} \\
 & + Q_{39} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\gamma}{\alpha_0+1}} + Q_{40} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}}
 \end{aligned}$$

for some constants  $Q_{36}, \dots, Q_{40}$ . This together with (2.17) and (2.18) implies that there is a constant  $Q_{41}$  with

$$(2.20) \quad \sup_{t \in [0,1]} |y(t)| \leq Q_{41}.$$

Now (2.19), (2.20) together with Theorem 1.2 establish the result. □

**Example.** Theorem 2.1 (here  $H_{\alpha_0, \theta}(u) = H_{\frac{1}{3}, \frac{1}{3}}(u)$ ) immediately guarantees that

$$\begin{cases}
 y'' + n^2\pi^2 y = y^{\frac{1}{3}} + [y']^{\frac{1}{7}} + 1 \quad \text{a.e. on } [0, 1] \\
 y(0) = y(1) = 0, \quad n \in \{1, 2, \dots\}
 \end{cases}$$

has a solution.

One can improve considerably the above theorem if  $m = 0$  (at the first eigenvalue). In particular the condition  $0 < \alpha_0 < 1$  is replaced by  $\alpha_0 > 0$  in this case; also condition (2.5) can be improved and the condition  $\sigma < \frac{\alpha_0}{2}$  can be relaxed. We present two existence results.

Consider

$$(2.21) \quad \begin{cases} \frac{1}{p}(py')' + ry + \lambda_0 qy = f(t, y, py') \quad \text{a.e. on } [0, 1] \\ y \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

where  $\lambda_0$  is the first eigenvalue of (2.2).

**Theorem 2.2.** *Let  $pf : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$  be an  $L^1$ -Carathéodory function with (1.3) and (1.9) satisfied. Suppose  $f$  has the decomposition  $f(t, u_1, u_2) = g(t, u_1, u_2) + h(t, u_1, u_2)$  with  $pg, ph : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$   $L^1$ -Carathéodory functions and*

$$(2.22) \quad \begin{cases} \text{there exist constants } A > 0, \alpha_0 > 0 \text{ and a function } \phi \in L^1_p[0, 1], \\ \phi > 0 \text{ a.e. on } [0, 1] \text{ with } u_1 g(t, u_1, u_2) \geq A\phi(t)H_{\alpha_0, \theta}(u_1) \\ \text{for a.e. } t \in [0, 1]; \text{ here } \alpha_0 \leq \theta \end{cases}$$

$$(2.23) \quad \begin{cases} \text{there exist } \phi_i \in L_p^1[0, 1], i = 1, 2, 3 \text{ and constants } \beta_0 \text{ and} \\ \sigma \text{ with } |h(t, u_1, u_2)| \leq \phi_1(t) + \phi_2(t)|u_1|^{\beta_0} + \phi_3(t)|u_2|^\sigma \text{ for} \\ \text{a.e. } t \in [0, 1]; \text{ here } \beta_0 < \alpha_0 \text{ and } \phi_3 > 0 \\ \text{a.e. on } [0, 1] \text{ or } \phi_3 \equiv 0 \text{ on } [0, 1] \end{cases}$$

$$(2.24) \quad \begin{cases} \text{there exist } \phi_i \in L_p^1[0, 1], i = 4, 5, 6 \text{ and constants } \gamma \leq \alpha_0, \tau > \sigma \\ \text{with } |g(t, u_1, u_2)| \leq \phi_4(t) + \phi_5(t)|u_1|^\gamma + \phi_6(t)|u_2|^\tau \\ \text{for a.e. } t \in [0, 1]; \\ \text{here } \phi_6 > 0 \text{ a.e. on } [0, 1] \text{ or } \phi_6 \equiv 0 \text{ on } [0, 1] \end{cases}$$

$$(2.25) \quad \sigma < \min\{1, \frac{\alpha_0}{\gamma}, \alpha_0\} \text{ and } \tau < 1$$

$$(2.26) \quad \begin{cases} \left(\phi_1^{\alpha_0+1}\phi^{-1}\right)^{\frac{1}{\alpha_0}} \in L_p^1[0, 1], \left(\phi_2^{\alpha_0+1}\phi^{-(\beta_0+1)}\right)^{\frac{1}{\alpha_0-\beta_0}} \in L_p^1[0, 1], \\ \left(\phi_5^{\alpha_0+1}\phi^{-\gamma}\right)^{\frac{1}{\alpha_0+1-\gamma}} \in L_p^1[0, 1] \text{ and } \left(\phi_3^{\alpha_0+1}\phi^{-1}\right)^{\frac{1}{\alpha_0}} \in L_p^1[0, 1] \end{cases}$$

and

$$(2.27) \quad \begin{cases} \text{with } \kappa = \max\{\frac{\alpha_0+1}{\alpha_0}, 2\}, \left(\phi_3\phi^{-\frac{1}{\alpha_0+1}}\right)^\kappa \in L_p^1[0, 1]. \text{ Also need} \\ \phi_6^\kappa \in L_p^1[0, 1] \text{ and } \left(\phi_3\phi^{-\frac{1}{\alpha_0+1}}\right)^{\frac{\kappa\tau}{\tau-\sigma}} (\phi_6)^{-\kappa\left(\frac{\sigma}{\tau-\sigma}\right)} \in L_p^1[0, 1] \\ \text{if } \phi_6 > 0 \text{ a.e. on } [0, 1] \end{cases}$$

holding. Then (2.21) has at least one solution  $y \in C[0, 1] \cap C^1(0, 1)$  with  $py' \in AC[0, 1]$ .

PROOF: Let  $y$  be a solution of (2.8) $_\lambda$  with  $m = 0$ . Following the ideas of Theorem 2.1 with  $u = 0$  and  $y = w$  we obtain the analogue of (2.11), namely

$$(2.28) \quad A \int_0^1 p\phi|y|^{\alpha_0+1} dt \leq A \int_0^1 p\phi dt + \int_0^1 p\phi_1|y| dt + \int_0^1 p\phi_2|y|^{\beta_0+1} dt + \int_0^1 p\phi_3|y||py'|^\sigma dt.$$

Hölder's inequality implies

$$(2.29) \quad A \int_0^1 p\phi|y|^{\alpha_0+1} dt \leq N_0 + N_1 \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{1}{\alpha_0+1}} + N_2 \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{\beta_0+1}{\alpha_0+1}} + \int_0^1 p\phi_3|y||py'|^\sigma dt$$

for some constants  $N_0, N_1$  and  $N_2$ . Let  $\kappa = \max\{2, \frac{\alpha_0+1}{\alpha_0}\}$ . Hölder's inequality together with assumption (2.27) implies

$$\int_0^1 p\phi_3|y||py'|^\sigma dt \leq \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{1}{\alpha_0+1}} \left(\int_0^1 p\left(\phi_3\phi^{-\frac{1}{\alpha_0+1}}\right)^\kappa |py'|^{\sigma\kappa} dt\right)^{\frac{1}{\kappa}}$$

if  $\kappa = \frac{\alpha_0+1}{\alpha_0}$  whereas

$$\int_0^1 p\phi_3|y||py'|^\sigma dt \leq \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{1}{\alpha_0+1}} \times \left(\int_0^1 p\left(\phi_3\phi^{-\frac{1}{\alpha_0+1}}\right)^\kappa |py'|^{\sigma\kappa} dt\right)^{\frac{1}{\kappa}} \left(\int_0^1 p(t) dt\right)^{\frac{\alpha_0-1}{2(\alpha_0+1)}}$$

if  $\kappa = 2$ . Put this into (2.29) and essentially the same reasoning as in Theorem 2.1 establishes the existence of constants  $N_3$  and  $N_4$  with

$$(2.30) \quad \int_0^1 p\phi|y|^{\alpha_0+1} dt \leq N_3 + N_4 \left(\int_0^1 p\left(\phi_3\phi^{-\frac{1}{\alpha_0+1}}\right)^\kappa |py'|^{\sigma\kappa} dt\right)^{\frac{\alpha_0+1}{\kappa\alpha_0}}.$$

Also (2.15) implies (as in Theorem 2.1) for  $t \in (0, 1)$  that

$$(2.31) \quad \begin{aligned} |p(t)y'(t)| \leq & N_5 + N_6 \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{\beta_0}{\alpha_0+1}} + N_7 \int_0^1 p\phi_3|py'|^\sigma dt \\ & + N_8 \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{\gamma}{\alpha_0+1}} + N_9 \int_0^1 p\phi_6|py'|^\tau dt \\ & + N_{10} \left(\int_0^1 p\phi|y|^{\alpha_0+1} dt\right)^{\frac{1}{\alpha_0+1}} \end{aligned}$$

for some constants  $N_5, \dots, N_{10}$ . Again with  $\kappa = \max\{2, \frac{\alpha_0+1}{\alpha_0}\}$  we have

$$\begin{aligned} \int_0^1 p\phi_3|py'|^\sigma dt & \leq \left(\int_0^1 p\left(\phi_3\phi^{-\frac{1}{\alpha_0+1}}\right)^\kappa |py'|^{\sigma\kappa} dt\right)^{\frac{1}{\kappa}} \left(\int_0^1 p\phi dt\right)^{\frac{1}{\alpha_0+1}} \\ & \quad \text{if } \kappa = \frac{\alpha_0+1}{\alpha_0} \\ \int_0^1 p\phi_3|py'|^\sigma dt & \leq \left(\int_0^1 p\left(\phi_3\phi^{-\frac{1}{\alpha_0+1}}\right)^\kappa |py'|^{\sigma\kappa} dt\right)^{\frac{1}{\kappa}} \left(\int_0^1 p\phi^{\frac{2}{\alpha_0+1}} dt\right)^{\frac{1}{2}} \\ & \quad \text{if } \kappa = 2 \\ \int_0^1 p\phi_6|py'|^\tau dt & \leq \left(\int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt\right)^{\frac{1}{\kappa}} \left(\int_0^1 p(t) dt\right)^{1-\frac{1}{\kappa}}. \end{aligned}$$

There are two cases to consider, namely  $\phi_6 > 0$  a.e. on  $[0, 1]$  or  $\phi_6 \equiv 0$  on  $[0, 1]$ .

Case (i).  $\phi_6 > 0$  a.e. on  $[0, 1]$ .

Putting the above into (2.31) and using (2.30) leads to

$$\begin{aligned}
 \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt &\leq N_{11} + N_{12} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^{\frac{\beta_0\tau}{\alpha_0}} \\
 &+ N_{13} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^\tau \\
 (2.32) \quad &+ N_{14} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^{\frac{\gamma\tau}{\alpha_0}} \\
 &+ N_{15} \left( \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \right)^\tau \\
 &+ N_{16} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^{\frac{\tau}{\alpha_0}}
 \end{aligned}$$

for some constants  $N_{11}, \dots, N_{16}$ . Also Hölder's inequality implies

$$\begin{aligned}
 &\int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \\
 &\leq \left( \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \right)^{\frac{\sigma}{\tau}} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^{\frac{\kappa\tau}{\tau-\sigma}} (\phi_6)^{-\kappa\left(\frac{\sigma}{\tau-\sigma}\right)} dt \right)^{\frac{\tau-\sigma}{\tau}}
 \end{aligned}$$

and putting this into (2.32) yields

$$\begin{aligned}
 \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt &\leq N_{17} + N_{18} \left( \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \right)^{\frac{\beta_0\sigma}{\alpha_0}} \\
 &+ N_{19} \left( \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \right)^\sigma + N_{20} \left( \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \right)^{\frac{\gamma\sigma}{\alpha_0}} \\
 &+ N_{21} \left( \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \right)^\tau + N_{22} \left( \int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \right)^{\frac{\sigma}{\alpha_0}}
 \end{aligned}$$

for some constants  $N_{17}, \dots, N_{22}$ . Now since  $\max\{\frac{\sigma\beta_0}{\alpha_0}, \sigma, \frac{\sigma\gamma}{\alpha_0}, \tau, \frac{\sigma}{\alpha_0}\} < 1$  then there exists a constant  $N_{23}$  with

$$\int_0^1 p\phi_6^\kappa |py'|^{\tau\kappa} dt \leq N_{23}$$

and essentially the same reasoning as in Theorem 2.1 establishes the result.



Case (ii).  $\phi_6 \equiv 0$  on  $[0, 1]$ .

We may assume without loss of generality that  $\sigma > 0$  and  $\phi_3 > 0$  a.e. on  $[0, 1]$ ; otherwise the result is easy. Then (2.31) for  $t \in (0, 1)$  becomes

$$\begin{aligned} |p(t)y'(t)| &\leq N_{23} + N_{24} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\beta_0}{\alpha_0+1}} \\ &\quad + N_{25} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^{\frac{1}{\kappa}} \\ &\quad + N_{26} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{\gamma}{\alpha_0+1}} + N_{27} \left( \int_0^1 p\phi|y|^{\alpha_0+1} dt \right)^{\frac{1}{\alpha_0+1}} \end{aligned}$$

for some constants  $N_{23}, \dots, N_{27}$ . This together with (2.30) leads to

$$\begin{aligned} &\int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \\ &\leq N_{28} + N_{29} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^{\frac{\beta_0\sigma}{\alpha_0}} \\ &\quad + N_{30} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^\sigma \\ &\quad + N_{31} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^{\frac{\gamma\sigma}{\alpha_0}} \\ &\quad + N_{32} \left( \int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \right)^{\frac{\sigma}{\alpha_0}} \end{aligned}$$

for some constants  $N_{28}, \dots, N_{32}$ . Thus there exists a constant  $N_{33}$  with

$$\int_0^1 p \left( \phi_3\phi^- - \frac{1}{\alpha_0+1} \right)^\kappa |py'|^{\sigma\kappa} dt \leq N_{33}$$

and the result follows as in Theorem 2.1. □

The next theorem establishes the existence of a nonnegative solution to

$$(2.33) \quad \begin{cases} \frac{1}{p}(py')' + \lambda_0 qy = \psi(t)f(t, y, py'), & 0 < t < 1 \\ y \in (\text{SL}) \text{ or } (\text{P}) \end{cases}$$

where  $\lambda_0$  is the first eigenvalue of (2.2) with  $r \equiv 0$  and  $q, \psi$  satisfies

$$(2.34) \quad q, \psi \in L^1_p[0, 1] \text{ with } q, \psi > 0 \text{ on } (0, 1).$$

Let

$$H^*_{\alpha_0, \theta}(u_1) = \begin{cases} u_1^{\theta+1}, & 0 \leq u_1 \leq 1 \\ u_1^{\alpha_0+1}, & 1 < u_1 < \infty. \end{cases}$$

**Theorem 2.3.** Let  $f : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$  be continuous with (1.3), (2.34) and

$$(2.35) \quad f(t, 0, 0) \leq 0$$

holding. Suppose  $\psi f$  has the decomposition  $\psi(t)f(t, u_1, u_2) = g(t, u_1, u_2) + h(t, u_1, u_2)$  with  $pg, ph : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$   $L^1$ -Carathéodory functions and

$$(2.36) \quad \begin{cases} \text{there exist constants } A > 0, \alpha_0 > 0 \text{ and a function } \phi \in L^1_p[0, 1], \\ \phi > 0 \text{ on } (0, 1) \text{ with } u_1g(t, u_1, u_2) \geq A\phi(t)H^*_{\alpha_0, \theta}(u_1) \\ \text{for } t \in (0, 1), u_1 \geq 0 \text{ and } u_2 \in \mathbf{R}; \text{ here } \alpha_0 \leq \theta \end{cases}$$

$$(2.37) \quad \begin{cases} \text{there exist } \phi_i \in L^1_p[0, 1], i = 1, 2, 3 \text{ and constants } \beta_0 \text{ and} \\ \sigma \text{ with } |h(t, u_1, u_2)| \leq \phi_1(t) + \phi_2(t)u_1^{\beta_0} + \phi_3(t)|u_2|^\sigma \text{ for} \\ t \in (0, 1), u_1 \geq 0 \text{ and } u_2 \in \mathbf{R}; \text{ here } \beta_0 < \alpha_0 \\ \text{and } \phi_3 > 0 \text{ on } (0, 1) \text{ or } \phi_3 \equiv 0 \end{cases}$$

and

$$(2.38) \quad \begin{cases} \text{there exist } \phi_i \in L^1_p[0, 1], i = 4, 5, 6 \text{ and constants } \gamma \leq \alpha_0, \tau > \sigma \\ \text{with } |g(t, u_1, u_2)| \leq \phi_4(t) + \phi_5(t)u_1^\gamma + \phi_6(t)|u_2|^\tau \text{ for } t \in (0, 1), u_1 \geq 0 \\ \text{and } u_2 \in \mathbf{R} \text{ and } \phi_6 > 0 \text{ on } (0, 1) \text{ or } \phi_6 \equiv 0 \end{cases}$$

hold. Finally suppose (2.25) and (2.26) are satisfied. Then (2.33) has at least one nonnegative solution  $y \in C[0, 1] \cap C^1(0, 1)$  with  $py' \in AC[0, 1]$ .

PROOF: Consider the family of problems

$$(2.39)_\lambda \quad \begin{cases} \frac{1}{p}(py')' + \mu qy = \lambda f^*(t, y, py'), \quad 0 < t < 1 \\ y \in (SL) \text{ or } (P) \end{cases}$$

where  $0 < \lambda < 1$  and

$$\mu = \begin{cases} 0 & \text{if } y \in (SL) \text{ and } \alpha^2 + a^2 > 0 \\ -1 & \text{if } y \in (P) \text{ or } y \in (SL) \text{ with } \alpha = a = 0. \end{cases}$$

Also

$$f^*(t, u_1, u_2) = \begin{cases} \psi(t)f(t, u_1, u_2) + (\mu - \lambda_0)qu_1, & u_1 \geq 0 \\ \psi(t)f(t, 0, u_2) + (\mu + 1)qu_1, & u_1 < 0. \end{cases}$$

Notice  $pf^* : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$  is an  $L^1$ -Carathéodory function.

Let  $y$  be a solution to (2.39) $_\lambda$  for some  $0 < \lambda < 1$ . We **claim** that  $y \geq 0$  on  $[0, 1]$ . If not then  $y$  would have a negative absolute minimum somewhere on  $[0, 1]$ ,

say at  $t_0$ . If  $t_0 \in (0, 1)$  then  $y'(t_0) = 0$  and this together with the differential equation and (2.35) yields

$$y''(t_0) = \frac{1}{p(t_0)}(p(t_0)y'(t_0))' = \lambda(\psi(t_0)f(t_0, 0, 0) + q(t_0)y(t_0)) + (\lambda - 1)\mu q(t_0)y(t_0) < 0,$$

a contradiction. Next suppose the negative absolute minimum were to occur at  $t_0 = 0$ . Consider first the Sturm Liouville boundary condition. Of course we need only consider  $\beta \neq 0$ . If  $\alpha \neq 0$  as well then

$$y(0) \lim_{t \rightarrow 0^+} p(t)y'(t) = \frac{\alpha}{\beta}y^2(0) > 0,$$

which implies  $y^2(t)$  is an increasing function near 0, a contradiction. So it remains to consider the case  $\alpha = 0$  and  $\beta \neq 0$ . The boundary condition is  $\lim_{t \rightarrow 0^+} p(t)y'(t) = 0$ . Now  $f(0, 0, 0) \leq 0$  and this together with the differential equation and (2.34) implies there exists  $\delta > 0$  with  $(p(t)y'(t))' < 0$  for  $t \in (0, \delta)$ . Thus the boundary condition implies  $p(t)y'(t) < 0$  for  $t \in (0, \delta)$ , a contradiction. Consequently  $t_0 \neq 0$ . A similar argument shows  $t_0 \neq 1$ . Thus our claim is established for Sturm Liouville boundary data.

Consider now Periodic boundary data. If the absolute minimum of  $y$  occurs at  $t_0 = 0$  then, since  $y(0) = y(1)$ , it must also occur at 1. Thus  $\lim_{t \rightarrow 0^+} p(t)y'(t) \geq 0$  and  $\lim_{t \rightarrow 1^-} p(t)y'(t) \leq 0$ . Consequently

$$\lim_{t \rightarrow 0^+} p(t)y'(t) = \lim_{t \rightarrow 1^-} p(t)y'(t) = 0$$

because of the second boundary condition. As above there exists  $\delta > 0$  with  $(p(t)y'(t))' < 0$  for  $t \in (0, \delta)$  and so  $p(t)y'(t) < 0$  for  $t \in (0, \delta)$ , a contradiction.

Thus  $y \geq 0$  on  $[0, 1]$  for any solution  $y$  to (2.39) $_{\lambda}$ . Consequently  $y$  satisfies

$$\frac{1}{p}(py')' + \mu qy = \lambda(\psi(t)f(t, y, py') + (\mu - \lambda_0)qy), \quad 0 < t < 1.$$

Essentially the same reasoning as in Theorem 2.2 (in this case we look at  $\int_0^1 p\phi y^{\alpha_0+1} dt$ ) guarantees the existence of a solution  $y$  to (2.39) $_1$ . Of course  $y$  is automatically a solution of (2.33) since  $y \geq 0$  on  $[0, 1]$ . □

**Existence theory II.**

In this subsection we examine the resonant problem (2.1) on the “right” of the eigenvalue.

**Theorem 2.4.** *Let  $pf : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$  be an  $L^1$ -Carathéodory function with (1.3) and (1.9) holding. Suppose  $f$  has the decomposition  $f(t, u_1, u_2) =$*

$g(t, u_1, u_2) + h(t, u_1, u_2)$  with  $pg, ph : [0, 1] \times \mathbf{R}^2 \rightarrow \mathbf{R}$   $L^1$ -Carathéodory functions and assume

$$(2.40) \quad \begin{cases} \text{there exist constants } A > 0, 0 < \alpha_0 < 1 \text{ and a function} \\ \phi \in L^1_p[0, 1], \phi > 0 \text{ a.e. on } [0, 1] \text{ with} \\ u_1g(t, u_1, u_2) \leq -A\phi(t)H_{\alpha_0, \theta}(u_1) \text{ for a.e. } t \in [0, 1]; \text{ here } \alpha_0 \leq \theta \end{cases}$$

holds. In addition assume (2.4), (2.5), (2.6) and (2.7) are satisfied. Then (2.1) has at least one solution  $y \in C[0, 1] \cap C^1(0, 1)$  with  $py' \in AC[0, 1]$ .

PROOF: Consider the family of problems

$$(2.41)_\lambda \quad \begin{cases} \frac{1}{p}(py')' + ry + \mu qy = \lambda[f(t, y, py') + (\mu - \lambda_m)qy] \text{ a.e. on } [0, 1] \\ y \in (SL) \text{ or } (P) \end{cases}$$

where  $0 < \lambda < 1$  and  $\lambda_m < \mu < \lambda_{m+1}$ .

Notice  $L^2_{pq}[0, 1] = \Gamma \oplus \Gamma^\perp$  where  $\Gamma = span \{\psi_0, \psi_1, \dots, \psi_m\}$ . Multiply (2.41) $_\lambda$  by  $w - u$  and integrate from 0 to 1 to obtain as in Theorem 2.1 ( $Q_0$  is as in Theorem 2.1)

$$(2.42) \quad \begin{aligned} & Q_0 + \int_0^1 [-p(w')^2 + prw^2 + \mu pqw^2] dt + \int_0^1 [p(u')^2 - pr u^2 - \mu pq u^2] dt \\ &= \lambda \int_0^1 (w - u)pf(t, y, py') dt + \lambda(\mu - \lambda_m) \int_0^1 pqw^2 dt \\ &\quad - \lambda(\mu - \lambda_m) \int_0^1 pq u^2 dt. \end{aligned}$$

Now since  $u \in \Gamma, w \in \Gamma^\perp$  and  $y = u + w$  we have

$$u = \sum_{i=0}^m c_i \psi_i \text{ and } w = \sum_{i=m+1}^\infty c_i \psi_i \text{ where } c_i = \langle y, \psi_i \rangle.$$

Also as before

$$\begin{aligned} & Q_0 + \int_0^1 [-p(w')^2 + prw^2 + \mu pqw^2] dt + \int_0^1 [p(u')^2 - pr u^2 - \mu pq u^2] dt \\ &\leq (\mu - \lambda_{m+1}) \int_0^1 pqw^2 dt + (\lambda_m - \mu) \int_0^1 pq u^2 dt \end{aligned}$$

so putting this into (2.42) yields

$$\begin{aligned} & \lambda \int_0^1 (w - u)pg(t, y, py') dt + (1 - \lambda)(\mu - \lambda_m) \int_0^1 pq u^2 dt \\ &\quad + (\lambda_{m+1} - \mu) \int_0^1 pqw^2 dt + \lambda(\mu - \lambda_m) \int_0^1 pqw^2 dt \\ &\leq -\lambda \int_0^1 (w - u)ph(t, y, py') dt. \end{aligned}$$

Now  $w - u = -y + 2w$  and  $-yg(t, y, py') \geq A\phi(t)H_{\alpha_0, \theta}(y)$  for a.e.  $t \in [0, 1]$  so with the above we have

$$\begin{aligned} A \int_0^1 p\phi H_{\alpha_0, \theta}(y) dt + (\mu - \lambda_m) \int_0^1 pqw^2 dt &\leq -2 \int_0^1 pwg(t, y, py') dt \\ &+ \int_0^1 p|y||h(t, y, py')| dt \\ &+ 2 \int_0^1 p|w||h(t, y, py')| dt. \end{aligned}$$

Essentially the same reasoning as in Theorem 2.1 (the only difference is that we use  $\int_0^1 pqw^2 dt$  in place of  $\int_0^1 pqu^2 dt$ ) establishes the result.  $\square$

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